# Vapor-liquid equilibrium (VLE) measurements of ethanol - heptane at isothermal ( $363.15,393.15$ and 423.15 K ) and isobaric ( 101.33 kPa ) conditions and correlation of liquid viscosity data 

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#### Abstract

Vapor-liquid equilibria (VLE) were measured for ethanol - heptane by a modified Rose-Williams equilibrium still combined with ebulliometer at three isotherms ( $363.15,393.15$ and 423.15 K ) and an isobar ( 101.33 kPa ). The azeotropes were observed in VLE, and the experimental data were correlated with the Peng-Robinson-Stryjek-Vera (PRSV) equation of state. The conventional mixing rule, with $k_{i j}$ equal to zero calculates the azeotropic pressure close to the vapor pressure of ethanol. The Adachi-Sugie (AS) mixing rule provides $0.282 \%$ average absolute relative deviation (AARD) in temperature for isobaric VLE and 1.711\% AARD for isothermal VLE. Eyring theory combined with the PRSV equation enabled an average of $2.810 \%$ AARD in correlation of literature liquid viscosity data, which is sufficient for considering biofuel injector. The method will be applied to flow analysis in general process design using mixtures because viscosity can be evaluated from equation of states.


## 1. Introduction

Electric vehicles (EVs) have received much attention as zero-local emission type of transportation. However, considering the present number of vehicles in the world, it is difficult to immediately switch from fuel-driven vehicles to EVs such that hybrid vehicles that use bioethanol blended with gasoline are being promoted in many countries. The volume percentage of ethanol in gasoline is given as an E-number, in which there are E5-10 grades in Europe, E10 and E15 grades in the U. S. and E18 to E27 grades in Brazil [1]. Bioethanol blended fuels are expected to be used in the future as fuels for vehicles. The physical properties, such as density, viscosity, heat capacity, and thermal conductivity, and phase equilibria are essential for design of biofuel injectors and their supply systems in gasoline engines [2-5]. Experimental vapor-liquid equilibria (VLE) data have been measured for ethanol and heptane mixtures by Katz and Newman [6] and Raal et al. [7]. Isothermal VLE and related data have been reported at 303 K [8, 9], at 298.15 K [10], 303 K [11], 303 K and $343 \mathrm{~K}[12], 313 \mathrm{~K}[13], 313 \mathrm{~K}[14,15], 343 \mathrm{~K}, 353 \mathrm{~K}$ and 363 K [16], $483 \mathrm{~K}, 508 \mathrm{~K}$ and 523 K [17]. However, there are few reports in the temperature range from 353 K to 453 K , which is necessary for the design of direct fuel injector systems [2,3]. The authors have reported VLE data using equilibrium still and ebulliometer not only for systems under isobaric conditions, but also isothermal conditions [18-

22]. In this research, experimental techniques were applied for determination of isothermal VLE data of the ethanol - heptane system at 363 K to 423 K . To supplement the isothermal VLE measurements in this work, isobaric measurements were also measured at 101.33 kPa to validate the reliability of the data. In mechanical engineering, REFPROP has been widely used for the estimation of physical properties like phase equilibria, viscosity, heat capacity and thermal conductivity. However it is not always sufficient for mixtures containing polar substances. Therefore experimental data were correlated with the equation of state proposed by Styrjek and Vera [23], because it is reliable for calculating saturated vapor pressures of both hydrocarbons and alcohols. In this research, the equation of state was also applied to the estimation of viscosity by using Eyring theory. Methods that use Eyring theory for estimating viscosity from VLE data have been proposed [24-27]. Ono et al. [28] also employed the Eyring theory for estimating viscosity from P-V-T data in the liquid phase. Estimated viscosities from VLE data of this work were compared with those reported by Brunson and Byers [29].

## 1. Experimental section

### 2.1. Materials

Table 1 lists the chemicals used in this research. Ethanol and heptane were dehydrated with molecular sieve 3A before the measurements and

[^0]Table 1
Chemicals used in this study.

|  | CAS No. | Supplier | ${ }^{\text {appurity }}$ |
| :--- | :--- | :--- | :--- |
| Ethanol | $64-17-5$ | Wako Pure Chemical Industries, Osaka, Japan | $>99.9 \mathrm{wt} \%$ |
| Heptane | $142-82-5$ | Wako Pure Chemical Industries, Osaka, Japan | $>99.9 \mathrm{wt} \%$ |

${ }^{\text {a }}$ Checked by gas chromatograph


Fig. 1. Schematic diagram of experimental apparatus; 1: Nitrogen cylinder; 2: Safety valve; 3: Solenoid valve; 4: Ebulliometer for reference / 4a: Pt resistance thermometer; 5: Surge tank; 6: Cold trap; 7: Pressure indicator; 8: Modified Rose-Williams equilibrium still for sample/ 8a: Pt resistance thermometer / 8b: Equilibrium chamber / 8c: Still / 8d,e: Sampling port; 9:Low temperature thermostatic bath.
purities were checked with a gas chromatograph. Peak areas of the impurities were no more than $0.1 \%$ for ethanol or heptane. Ion exchanged distilled water was used as a reference to determine experimental pressures at isothermal conditions.

### 2.2. Equipment and procedure

Fig. 1 shows the apparatus employed in this research with experimental procedures described in detail in previous reports [22]. The apparatus was composed of a modified Rose-Williams still, an ebulliometer for reference (part 4 in Fig. 1), a surge tank (part 5), a nitrogen cylinder (part 1), a low temperature thermostatic bath (part 9) and other equipment with the main parts of the apparatus being made of stainless steel 304. The apparatus was designed to have a safe working pressure of 5 MPa . The sample was heated up to about 483 K by a cartridge heater installed in the equilibrium still. In the procedure, the sample solution was loaded into the still of the ebulliometer (part 8c), and boiled under a given pressure. Then, the vapor and the liquid of the sample were sent together to the equilibrium chamber (part 8 b ). The liquid was returned to the still via the sampling port (part 8d). Otherwise, the vapor was sent to a Liebig condenser and liquefied. The liquefied vapor was returned to the still together with the liquid from the equilibrium chamber. In the apparatus operation, the number of drops from the condenser were ensured to be constant implying that the system was at steadystate and that the system was in the equilibrium. In isothermal VLE measurements, temperature was measured with a platinum resistance thermometer (part 8a), and pressure was evaluated from the temperature of pure water measured by another thermometer (part 4a). The two
thermometers were calibrated at the National Institute of Advanced Industrial Science and Technology (AIST), Japan. An equation, proposed by Bridgeman and Aldrich [30], was employed for evaluation of pressure balance with that of the ebulliometer for the sample. A small portion of the liquid phase and of the liquefied vapor phase were taken from sampling ports (parts 8d and 8e) using a gastight syringe, and the composition was analyzed with a gas chromatograph (Shimadzu GC-4BT, Kyoto, Japan) equipped with a thermal conductivity detector. The column packing was PEG-1500, 15\% polyethylene glycol on Uniport, 60/80 mesh (Gasukuro Kogyo, Tokyo). In the isobaric VLE measurements, the system was at atmospheric pressure and the temperature was measured solely by the thermometer (part 8a), while the ebulliometer was not used.

In the isobaric measurements, the temperature measured was corrected to that at 101.33 kPa via:
$T=T_{\text {meas }}+\frac{1}{(\ln 10) \sum_{i} \frac{x_{i} B_{i}}{\left(T_{\text {meas }}+C_{i}\right)^{2}}} \times \frac{101.33-p_{\text {atm, meas }} / k P a}{p_{\text {atm,meas }} / k P a}$
where $T_{\text {meas }}$ is the temperature measured under atmospheric pressure, $p_{\text {atm,meas }}, x_{i}$ is the mole fraction of liquid phase. $B_{i}$ and $C_{i}$ are constants in the following Antione equation:
$\ln p^{s}{ }_{i}=A_{i}-\frac{B_{i}}{T / K-C_{i}}$
where constants, $A_{i}, B_{i}$ and $C_{i}$, were determined from the experimental saturated vapor pressure data of ethanol and heptane. Considering that the thermometer calibration and thermal oscillation of the system, uncertainty of temperature was estimated to be $u(T)=0.06 \mathrm{~K}$.

Table 2
Critical properties used for PRSV equation [23].

|  | Critical temperature $T_{c} / \mathrm{K}$ | Critical pressure $p_{c} / \mathrm{kPa}$ | Acentric factor $\omega /-$ | $\kappa_{1} /-$ |
| :--- | :--- | :--- | :--- | :--- |
| Ethanol | 513.92 | 6148 | 0.64439 | -0.03374 |
| Heptane | 540.10 | 2735.75 | 0.35022 | 0.04648 |

Considering the peak area in gas chromatogram, uncertainties in mole fraction for vapor and liquid were estimated to be $u\left(x_{i}\right)=u\left(y_{i}\right)=0.003$, respectively. The uncertainty of pressure depended on the temperature of water in the ebulliometer, because the value of pressure was determined by the equation of Bridgeman and Aldrich [30]. For example, the uncertainties of pressure at (100, 500, 1000, 2000) kPa were evaluated to be $u(p)=(0.25,0.70,1.20,2.43) \mathrm{kPa}$, respectively.

## 3. Correlation

### 3.1. VLE correlation

The modified Peng-Robinson equation of state of Stryjek and Vera (PRSV) [23] was employed for correlation of VLE data:
$p=\frac{R T}{v-b}-\frac{a}{v^{2}+2 b v-b^{2}}$
where parameters, $a$ and $b$, were evaluated from the critical temperature, $T_{c}$, critical pressure, $p_{c}$, acentric factor, $\omega$, and an additional parameter, $\kappa_{1}$ [23]. The additional parameter $\kappa_{1}$ was proposed for well representing the saturated vapor pressure data [23]. Considering the azeotrope reported in ethanol -heptane [8-17], the good reproducibility is essential for saturated vapor pressure to determine the binary parameters in the mixing rule. Table 2 lists the parameters used in the calculations. Some mixing rules should be considered for the correlation of VLE data, because van der Waals one fluid model does not always provide azeotropes. It is well known that mixing rules, based on excess free energy models [31-33], are employed for correlating the systems with complex phase behavior. However, the estimation of VLE and viscosity are considered for design of biofuel injectors in gasoline engine. Then, complex mathematical functions are not useful for estimation of heat capacities. Though estimation of heat capacity were not discussed in this research, mixing rules with simple mathematical function will be preferable. Therefore, mixing rules proposed by Adachi and Sugie (AS) were employed, since these also allow description of polar-nonpolar systems [34]:
$a=\sum_{i} \sum_{j} x_{i} x_{j}\left[1-k_{i j}+l_{i j}\left(x_{i}-x_{j}\right)\right]\left(a_{i} a_{j}\right)^{1 / 2}$
$b=\sum_{i} \sum_{j} x_{i} x_{j} \frac{b_{i}+b_{j}}{2}$
where $k_{i j}$ and $l_{i j}$ are two binary parameters. It should be noted that the functional form of Eq. (4) is similar to that of a Margules type activity coefficient model. In this research, the binary parameters were determined from fitting isothermal VLE data. Several combinations of ( $k_{i j}$, $l_{i j}$ ) were obtained to provide the azeotropic point at first and successively, with final values being determined by minimizing the following objective function:
O.F. $=\sum\left(x_{1}-x_{1, \text { calc }}\right)$

After identifying the temperature dependence in $k_{i j}$ and $l_{i j}$, the isobaric VLE were calculated.

### 3.2. Estimation of liquid viscosity

According to Eyring theory [35], the liquid viscosity of the binary mixture, $\eta$, at given temperature, $T$, and pressure, $p$, is given as follows:
$\ln \eta v=x_{1} \ln \left[\eta_{1}^{\circ}(T, p) v_{1}^{\circ}(T, p)\right]+x_{2} \ln \left[\eta_{2}^{\circ}(T, p) v_{2}^{\circ}(T, p)\right]+\frac{E_{a}}{R T}$
where $v_{i}^{\circ}$ is the molar volume, $\eta_{i}^{\circ}$ is the viscosity of pure component $i$ and $E_{a}$ is the activation energy that can be assumed to be proportional to the excess Gibbs energy, $G^{E}$ :
$E_{a}=\sigma G^{E}$
where $\sigma$ is a constant that depends on the system and $G^{E}$ is defined as follows:
$\frac{G^{E}}{R T}=x_{1} \ln \frac{f_{1}(T, p)}{x_{1} f_{1}^{\circ}(T, p)}+x_{2} \ln \frac{f_{2}(T, p)}{x_{2} f_{2}^{\circ}(T, p)}$
where $f_{i}^{\circ}$ and $f_{i}$ are the fugacities of component $i$ as a pure substance and as a component in the mixture, respectively. In this research, $v, v_{i}^{\circ}, f_{i}^{\circ}$ and $f_{i}$ were estimated with the PRSV equation.

### 3.3. Deviations and absolute relative deviations

The relative deviations (RDs) and the absolute relative deviations (ARDs) were evaluated as follows:

$$
\begin{equation*}
\frac{\Delta X}{X_{\text {exp }}}=\frac{X-X_{\text {exp }}}{X_{\text {exp }}} \tag{RD}
\end{equation*}
$$

(ARD) $\quad \frac{|\Delta X|}{X_{\text {exp }}}=\frac{\left|X-X_{\text {exp }}\right|}{X_{\text {exp }}}$
Sparacwhere $X$ and $X_{\text {exp }}$ refer to calculated and experimental properties. For saturated pressures, $X$ corresponds to that of the Antoine equation $p_{\text {fit }}^{s}$ or that of the PRSV equation $p_{c a l}^{s}$ at a given temperature. For isothermal VLE, $X$ is the pressure given by the PRSV equation at a given mole fraction $x_{1}$ and mole fraction $y_{1}$ at given pressure $p$. For isothermal VLE, $X$ is the temperature at a given mole fraction $x_{1}$ and mole fraction $y_{1}$ at a given temperature $T$. For the estimation of viscosity, $X$ is the viscosity $\eta$ at a given mole fraction $x_{1}$.

## 4. Result and discussion

### 4.1. Saturated vapor pressure

Table 3 gives the saturated vapor pressures for ethanol and heptane. The experimental pressure should be higher than atmospheric pressure in the apparatus employed, so the experimental temperatures ranged from ( 351.48 to 463.08 ) K and ( 371.82 to 483.15 ) K for ethanol and heptane, respectively. The data for ethanol agreed with the standard data of Schroder et al. [36], and those for heptane agreed with values reported by National Institute of Standard and Technology (NIST) [37]. The constants in the Antoine equation, Eq. (2), were determined by fitting experimental vapor pressure data (Table 4). The RDs are shown in Fig. 2 and the average absolute relative deviation (AARDs) were $0.295 \%$ and $0.270 \%$ for the experimental data of ethanol and heptane, respectively. The AARDs were also calculated for the literature data ethanol and the standard data of heptane. The AARD for the literature data of ethanol were $0.471 \%$. The AARD for the standard data of heptane was calculated for every 0.15 K from 371.85 K to 383.20 K in the NIST Chemistry WebBook, and the AARD was $0.180 \%$. Constants in the Antoine equation were used in the isobaric VLE to correct the pressure from atmospheric pressure at measurement to 101.33 kPa .

Table 3
Saturated vapor pressure for ethanol(1) and heptane(2).

| Measured |  | Antoine equation |  |  | PRSV equation |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{\text {a }}$ Temperature | ${ }^{\text {b }}$ Pressure | Pressure |  |  | Pressure |  |  |
| T/K | $p^{s} / \mathrm{kPa}$ | $p_{\text {fit }}^{s} /-$ | ${ }^{\text {c }} \Delta p^{s} /-$ | ${ }^{\mathrm{d}}\left\|\Delta p^{s}\right\| / p^{s}$ exp $/ \%$ | $p_{\text {cal }}{ }^{\text {/- }}$ | ${ }^{\text {e }} \Delta p^{s} /-$ | ${ }^{\mathrm{f}}\left\|\Delta p^{s}\right\| / p^{s}$ exp $/ \%$ |
| Ethanol (1) |  |  |  |  |  |  |  |
| 351.48 | 101.77 | 101.61 | -0.16 | 0.160 | 101.77 | 0.00 | 0.000 |
| 363.15 | 158.48 | 158.26 | -0.22 | 0.140 | 158.00 | -0.48 | 0.303 |
| 373.15 | 225.21 | 225.24 | 0.03 | 0.013 | 224.55 | -0.66 | 0.293 |
| 383.15 | 312.45 | 313.52 | 1.07 | 0.341 | 312.04 | -0.41 | 0.131 |
| 393.22 | 426.23 | 428.57 | 2.34 | 0.548 | 425.78 | -0.45 | 0.106 |
| 403.15 | 571.38 | 572.73 | 1.35 | 0.236 | 568.02 | -3.36 | 0.588 |
| 413.15 | 755.33 | 754.21 | -1.12 | 0.148 | 746.81 | -8.52 | 1.128 |
| 423.16 | 983.30 | 978.27 | -5.03 | 0.512 | 967.38 | -15.92 | 1.619 |
| 433.20 | 1257.89 | 1251.92 | -5.97 | 0.475 | 1236.86 | -21.03 | 1.672 |
| 443.14 | 1582.62 | 1577.65 | -4.97 | 0.314 | 1558.13 | -24.49 | 1.547 |
| 453.14 | 1966.01 | 1967.21 | 1.20 | 0.061 | 1943.47 | -22.54 | 1.146 |
| 463.08 | 2408.79 | 2422.90 | 14.11 | 0.586 | 2396.22 | -12.57 | 0.522 |
|  |  | average | 0.22 | 0.295 | average | -9.20 | 0.755 |
| Heptane (2) |  |  |  |  |  |  |  |
| 371.82 | 102.53 | 102.97 | 0.44 | 0.431 | 101.92 | -0.61 | 0.595 |
| 382.99 | 140.68 | 140.57 | -0.11 | 0.077 | 139.95 | -0.73 | 0.519 |
| 393.16 | 184.39 | 183.58 | -0.81 | 0.438 | 183.22 | -1.17 | 0.635 |
| 403.10 | 236.03 | 235.02 | -1.01 | 0.428 | 234.94 | -1.09 | 0.462 |
| 413.22 | 299.21 | 298.31 | -0.90 | 0.299 | 298.51 | -0.70 | 0.234 |
| 423.15 | 371.21 | 372.54 | 1.33 | 0.359 | 373.00 | 1.79 | 0.482 |
| 433.15 | 459.10 | 460.90 | 1.80 | 0.391 | 461.63 | 2.53 | 0.551 |
| 443.15 | 562.77 | 564.43 | 1.66 | 0.294 | 565.48 | 2.71 | 0.482 |
| 453.16 | 683.67 | 684.82 | 1.15 | 0.168 | 686.34 | 2.67 | 0.391 |
| 463.15 | 823.31 | 823.28 | -0.03 | 0.003 | 825.53 | 2.22 | 0.270 |
| 473.15 | 983.16 | 981.81 | -1.35 | 0.138 | 985.27 | 2.11 | 0.215 |
| 483.15 | 1164.39 | 1161.88 | -2.51 | 0.215 | 1167.34 | 2.95 | 0.253 |
|  |  | average | -0.03 | 0.270 | average | 1.06 | 0.424 |

${ }^{\text {a }} u(T)=0.06 \mathrm{~K}$
${ }^{\mathrm{b}} u\left(p^{s}\right)$ depends on the pressure range, 0.25 kPa (at 100.00 kPa ), 0.70 kPa (at 500.00 kPa ), 1.20 kPa (at 1000.00 kPa ) and 2.43 kPa (at 2000.00 kPa )
${ }^{\mathrm{c}} \Delta p^{s}=p_{\text {fit }}^{s}-p_{\text {exp }}^{s}$
${ }^{\mathrm{d}}\left|\Delta p^{s}\right| / p_{\text {exp }}^{s}=\frac{\left|p_{f i t}-p_{\text {exp }}^{s}\right|}{p_{\text {exp }}^{s}} \times 100$
${ }^{\mathrm{e}} \Delta p^{s}=p_{\text {cal }}^{s}-p_{\text {exp }}^{s}$
${ }^{\mathrm{f}}\left|\Delta p^{s}\right| / p_{\text {exp }}^{s}=\frac{\left|p_{c a l}^{s}-p_{\text {exp }}^{s}\right|}{p_{\text {exp }}^{s}} \times 100$

Table 4
Constants in Antoine equation determined from data in Table 3.

|  | Constants in ${ }^{\mathrm{a}}$ Antoine equation |  | Temperature range |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- |
|  | $A_{i} /-$ | $B_{i} / \mathrm{K}$ | $C_{i} / \mathrm{K}$ | $T_{\min } / \mathrm{K}$ | $T_{\max } / \mathrm{K}$ |
| Ethanol (1) | 15.9114 | 3225.34 | -65.81 | 351.48 | 463.08 |
| Heptane (2) | 14.5964 | 3450.11 | -25.49 | 371.82 | 483.15 |
| ${ }^{\mathrm{a}} \ln p^{s}{ }_{i} / k P a=A_{i}-\frac{B_{i}}{T / K+C_{i}}$ |  |  |  |  |  |

### 4.2. Isobaric VLE

Table 5 lists experimental data of VLE for ethanol (1) - heptane (2) at 101.33 kPa . The thermodynamic consistency test, proposed by Ferdenslund et al. [38, 39], was carried out for the experimental data, and its reliability was ensured. Detailed descriptions are given in Supplemental Materials 1. Fig. 3 shows isobaric VLE along with the data of Katz and Newman [6], and Raal et al. [7] at 760 mmHg ( 101 kPa ). The experimental temperatures rapidly decreased in mole fraction of ethanol up to $x_{1}=0.150$. Though Katz and Newman [6] reported constant temperatures in the mole fraction of ethanol from 0.489 to 0.791 , the experimental data and those of Raal et al. [7] showed a lower temperature. The VLE data were considered to be showing an azeotrope in this work. The azeotropic point is sometimes affected by a small amount of water in ethanol, and can changed into a heterogeneous azeotrope.


Fig. 2. RDs for saturated vapor pressure of ethanol and heptane; ( $\bigcirc$ ): ethanol for Antoine equation; $(\bullet)$ : heptane for Antoine equation; $(\triangle)$ : ethanol for PRSV equation; ( $\mathbf{\Delta}$ ): heptane for PRSV equation; ( $\boldsymbol{\square}$ ): ethanol (Schroder et al. [33]) for Antoine equation; (--) : heptane (NIST Chemistry WebBook [34]) for Antoine equation.

Though it is difficult to check heterogeneous phases in the present apparatus, cloudy solutions indicative of heterogeneous azeotropy were not observed in the measurements. Using the experimental data, the azeotropic point was determined by plots of $y_{1}-x_{1}$, and an equation for the temperature by using Lagrange interpolation. The azeotropic point determined is listed in Table 6.

Table 5
Isobaric VLE for ethanol (1) - heptane (2) at ${ }^{\mathrm{a}} 101.33 \mathrm{kPa}$.

| Measured |  |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |

${ }^{\text {a }} u(p)=0.25 \mathrm{kPa}$
${ }^{\mathrm{b}} u(T)=0.06 \mathrm{~K}$
c $u\left(x_{1}\right)=u\left(y_{1}\right)=0.003$
${ }^{\mathrm{d}} \Delta T=T_{\text {cal }}-T_{\exp }$
${ }^{\mathrm{e}}|\Delta T| / T_{\text {exp }}=\frac{\left|T_{\text {cal }}-T_{\text {exp }}\right|}{T_{\text {exp }}} \times 100$
${ }^{\mathrm{f}} \Delta y_{1}=y_{1, \text { cal }}-y_{1, \text { exp }}$
$\mathrm{g}\left|\Delta y_{1}\right| / y_{1, \exp }=\frac{\left|y_{1, \text { cal }}-y_{1, e x p}\right|}{y_{1, \text { exp }}} \times 100$

Table 6
Measured and calculated azeotropes for ethanol(1) - heptane(2).

| Type of data | Measured <br> Mole fraction of ethanol $x_{1}, a z /-$ | Temperature $T_{a z} / \mathrm{K}$ | Pressure $p_{a z} / \mathrm{kPa}$ | PRSV equation <br> Mole fraction of ethanol $x_{1}, a z /-$ | Temperature $T_{a z} / \mathrm{K}$ | Pressure $p_{a z} / \mathrm{kPa}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Isobaric | 0.645 | 344.48 | ${ }^{\text {a }} 101.33$ | 0.651 | 344.87 | ${ }^{\text {a }} 101.33$ |
| Isothermal | 0.715 | ${ }^{\text {a }} 363.15$ | 199.55 | 0.680 | ${ }^{\text {a }} 363.15$ | 198.65 |
| Isothermal | 0.749 | ${ }^{\text {a }} 393.15$ | 497.41 | 0.734 | ${ }^{\text {a }} 393.15$ | 494.23 |
| Isothermal | 0.781 | ${ }^{\text {a }} 423.15$ | 1066.47 | 0.745 | ${ }^{\text {a }} 423.15$ | 1069.17 |

${ }^{\text {a }}$ fixed at measurement

### 4.3. Isothermal VLE

Table 7 lists experimental VLE data for ethanol (1) - heptane (2) at $(363.15,393.15,423.15)$ K. Fig. 4 shows isothermal VLE along with data reported by Ramalho and Delmas [16] at 343 K and 363 K , Berro et al. [12] at 343.17 K , and Seo et al. [17] at 483.15 K . The VLE data of the five isotherms had similar tendencies, however, the experimental pressures at 363.15 K were slightly lower than those of Ramalho and Delmas [16] at 363 K . The pressures in the data of Berro et al. [12] were also lower than those of Ramalho and Delmas [16] at 343 K . Ramalho and Delmas [16], and Berro et al. [12] measured VLE with an equilibrium still, while Seo et al. [17] made measurements with a high-pressure cell equipped with circulation pumps. When heat insulation of the equilibrium chamber is not sufficient in the equilibrium still, then the measured temperature probably tends to be slightly high, however, detailed discussion on this point is not possible here. The azeotropic points are shown in Table 6. The mole fraction of ethanol at the azeotropic point shifted to be larger with increasing temperature (Table 6). According to Seo et al. [17], the azeotropic point is separated by two points at high temperatures.

### 4.4. Calculation of saturated vapor pressure with PRSV equation

Prior to VLE correlation, reproducibility of the PRSV equation was checked for reliability of calculating experimental saturated vapor pres-
sure data. Table 3 lists the calculation results and Fig. 2 shows the RDs, which were all negative for ethanol over the entire range of temperatures. The largest RD in PRSV calculation of vapor pressure was found around 433 K . Even though the RDs for PRSV were larger than those for the Antoine equation, the AARD (Table 3) was $0.755 \%$. The sign of the RDs for PRSV changed from negative to positive for heptane, with deviatios being smaller than those for ethanol and having the AARD (Table 3) of $0.424 \%$.

### 4.5. Correlation of VLE

At first, VLE at 393.15 K was correlated with conventional mixing rules setting $l_{12}=0.000$ (Supplemental Materials 2). For $k_{12}=0.000$, the azeotropic point was calculated to be at a mole fraction close to ethanol, otherwise, $k_{12}=0.170$ provided reproducibility only at a mole fraction close to heptane, suggesting that at least two parameters are necessary in the mixing rules. It is well known that binary parameters, $k_{12}$ and $l_{12}$, can have temperature dependence, so their values were determined separately by fitting the data at three temperature (Table 8). Table 7 and Fig. 4 compare calculated results with the data and Table 6 shows calculated azeotropic points with the PRSV equation. Fig. 5 shows the RDs for pressures and vapor phase mole fractions with the correlation mainly being carried out to reproduce the azeotropic point. An enlargement of composition ranges around the azeotrope is shown in Fig. 5 for which the AARDs (Table 7) for pressure were $2.457 \%, 1.392 \%$ and $1.340 \%$

Table 7
Isothermal VLE for ethanol (1) - heptane (2) at 363.15, 393.15 and 423.15 K .

| Measured |  |  | PRSV equation at $x_{1, \exp }$ |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| ${ }^{a}$ Pressure | ${ }^{\text {b }}$ Mole fraction of ethanol |  | Pressure |  |  | Mole fraction of ethanol in vapor |  |  |
| $p / \mathrm{kPa}$ | $x_{1} /$ - | $y_{1} /-$ | $p_{\text {cal }} / \mathrm{kPa}$ | ${ }^{\mathrm{c}} \Delta p / \mathrm{kPa}$ | ${ }^{\mathrm{d}}\|\Delta p\| / p_{\text {exp }} / \%$ | $y_{1, \text { cal }} /-$ | ${ }^{\mathrm{e}} \Delta y_{1} /-$ | ${ }^{\mathrm{f}}\left\|\Delta y_{1}\right\| / y_{1, \exp } / \%$ |
| ${ }^{8} 363.15$ K |  |  |  |  |  |  |  |  |
| 152.58 | 0.070 | 0.500 | 133.18 | -19.40 | 12.712 | 0.434 | -0.066 | 13.143 |
| 173.62 | 0.122 | 0.580 | 157.95 | -15.67 | 9.023 | 0.536 | -0.044 | 7.531 |
| 181.52 | 0.171 | 0.608 | 173.15 | -8.37 | 4.608 | 0.587 | -0.021 | 3.450 |
| 185.92 | 0.213 | 0.628 | 181.84 | -4.08 | 2.193 | 0.614 | -0.014 | 2.267 |
| 189.26 | 0.277 | 0.648 | 189.97 | 0.71 | 0.376 | 0.637 | -0.011 | 1.755 |
| 193.41 | 0.368 | 0.667 | 195.35 | 1.94 | 1.004 | 0.657 | -0.010 | 1.567 |
| 194.97 | 0.458 | 0.680 | 197.32 | 2.35 | 1.207 | 0.665 | -0.015 | 2.223 |
| 197.47 | 0.533 | 0.692 | 197.94 | 0.47 | 0.240 | 0.669 | -0.023 | 3.314 |
| 198.38 | 0.615 | 0.701 | 198.48 | 0.10 | 0.050 | 0.674 | -0.027 | 3.824 |
| 199.44 | 0.686 | 0.708 | 198.65 | -0.79 | 0.396 | 0.679 | -0.029 | 4.086 |
| 199.43 | 0.750 | 0.719 | 198.25 | -1.18 | 0.590 | 0.694 | -0.025 | 3.505 |
| 198.66 | 0.805 | 0.730 | 196.84 | -1.82 | 0.918 | 0.712 | -0.018 | 2.463 |
| 191.98 | 0.879 | 0.782 | 191.53 | -0.45 | 0.236 | 0.757 | -0.025 | 3.153 |
| 184.69 | 0.930 | 0.832 | 183.13 | -1.56 | 0.846 | 0.817 | -0.015 | 1.781 |
|  |  |  | average | -3.41 | 2.457 | average | -0.024 | 3.862 |
| ${ }^{8} 393.15$ K |  |  |  |  |  |  |  |  |
| 249.31 | 0.028 | 0.273 | 237.08 | -12.23 | 4.905 | 0.232 | -0.041 | 14.970 |
| 421.18 | 0.182 | 0.614 | 405.08 | -16.10 | 3.822 | 0.577 | -0.037 | 5.966 |
| 449.30 | 0.289 | 0.637 | 450.87 | 1.57 | 0.350 | 0.637 | -0.036 | 5.291 |
| 465.67 | 0.367 | 0.702 | 468.30 | 2.63 | 0.565 | 0.659 | -0.043 | 6.068 |
| 474.57 | 0.448 | 0.713 | 479.23 | 4.66 | 0.982 | 0.676 | -0.037 | 5.135 |
| 485.92 | 0.542 | 0.726 | 487.08 | 1.16 | 0.239 | 0.693 | -0.033 | 4.602 |
| 490.02 | 0.608 | 0.732 | 490.82 | 0.80 | 0.163 | 0.704 | -0.028 | 3.804 |
| 494.87 | 0.672 | 0.736 | 493.28 | -1.59 | 0.322 | 0.717 | -0.019 | 2.544 |
| 497.08 | 0.721 | 0.742 | 494.18 | -2.90 | 0.583 | 0.730 | -0.012 | 1.643 |
| 497.16 | 0.781 | 0.759 | 493.43 | -3.73 | 0.750 | 0.750 | -0.009 | 1.159 |
| 495.60 | 0.837 | 0.785 | 489.51 | -6.09 | 1.228 | 0.778 | -0.007 | 0.938 |
| 481.14 | 0.920 | 0.845 | 472.05 | -9.09 | 1.888 | 0.848 | 0.003 | 0.317 |
| 475.12 | 0.940 | 0.870 | 464.18 | -10.94 | 2.302 | 0.874 | 0.004 | 0.467 |
|  |  |  | average | -3.99 | 1.392 | average | -0.023 | 4.070 |
| 8423.15 K |  |  |  |  |  |  |  |  |
| 887.61 | 0.135 | 0.602 | 798.78 | -88.85 | 10.010 | 0.526 | -0.076 | 12.597 |
| 921.08 | 0.194 | 0.642 | 885.23 | -35.85 | 3.892 | 0.577 | -0.065 | 10.162 |
| 942.85 | 0.242 | 0.666 | 931.70 | -11.15 | 1.182 | 0.602 | -0.064 | 9.595 |
| 953.92 | 0.277 | 0.679 | 956.59 | 2.67 | 0.280 | 0.616 | -0.063 | 9.323 |
| 970.79 | 0.316 | 0.696 | 978.12 | 7.33 | 0.756 | 0.628 | -0.068 | 9.790 |
| 989.43 | 0.373 | 0.714 | 1001.70 | 12.27 | 1.240 | 0.642 | -0.072 | 10.051 |
| 1012.13 | 0.450 | 0.729 | 1024.78 | 12.65 | 1.250 | 0.659 | -0.070 | 9.637 |
| 1042.13 | 0.547 | 0.750 | 1046.33 | 4.20 | 0.403 | 0.680 | -0.070 | 9.368 |
| 1050.05 | 0.601 | 0.753 | 1055.81 | 5.76 | 0.548 | 0.693 | -0.060 | 7.924 |
| 1055.93 | 0.652 | 0.760 | 1062.95 | 7.02 | 0.665 | 0.708 | -0.052 | 6.778 |
| 1062.24 | 0.717 | 0.769 | 1068.49 | 6.25 | 0.589 | 0.733 | -0.036 | 4.725 |
| 1065.75 | 0.757 | 0.775 | 1069.05 | 3.30 | 0.310 | 0.740 | -0.035 | 4.459 |
| 1065.65 | 0.812 | 0.789 | 1064.67 | -0.98 | 0.092 | 0.783 | -0.006 | 0.718 |
| 1062.67 | 0.843 | 0.804 | 1058.68 | -3.99 | 0.375 | 0.806 | 0.002 | 0.217 |
| 1052.09 | 0.882 | 0.834 | 1046.57 | -5.52 | 0.524 | 0.840 | 0.006 | 0.714 |
| 1033.03 | 0.920 | 0.873 | 1028.80 | -4.23 | 0.410 | 0.879 | 0.006 | 0.713 |
| 1019.90 | 0.939 | 0.897 | 1017.34 | -2.56 | 0.251 | 0.904 | 0.007 | 0.755 |
|  |  |  | average | -5.39 | 1.340 | average | -0.042 | 6.325 |

${ }^{\text {a }} u(p)$ depends on the pressure range, 0.70 kPa (at 500.00 kPa ) and 1.20 kPa (at 1000.00 kPa )
${ }^{\mathrm{b}} u\left(x_{1}\right)=u\left(y_{1}\right)=0.003$
${ }^{\text {c }} \Delta p=p_{\text {cal }}-p_{\text {exp }}$
${ }^{\mathrm{d}}|\Delta p| / p_{\text {exp }}=\frac{\left|p_{\text {cat }}-p_{\text {exp }}\right|}{p_{\text {exp }}} \times 100$
e $\Delta y_{1}=y_{1, \text { cal }}-y_{1, \exp }$
f $\left|\Delta y_{1}\right| / y_{1, \text { exp }}=\frac{\left|y_{1, c a t}-y_{1, \text { exp }}\right|}{y_{1, e x}} \times 100$
${ }^{\mathrm{g}} u(T)=0.06 \mathrm{~K}$
for $363.15 \mathrm{~K}, 393.15 \mathrm{~K}$ and 423.15 K , respectively, while those for mole fraction in the vapor phase were $3.862 \%, 4.070 \%$ and $6.325 \%$, respectively. Especially, RDs were larger at small mole fractions of ethanol in the liquid phase. Calculations were carried out for the data of Berro et al. [12] at 343.15 K and Seo et al. [17] at 483.15 K . The fitted values of $k_{12}$ and $l_{12}$ are listed in Table 8, and the calculation results are shown in Fig. 4. The calculation results agreed well the experimental data for mole fractions in liquid phase at 483.15 K (Fig. 4). Otherwise, the cal-
culated mole fractions in the vapor phase (Fig. 4) had higher deviations compared with those in liquid phase. The temperature dependence of $k_{12}$ and $l_{12}$ were determined to be as follows:
$k_{12}=1.3193-1.1204 \times 10^{-3} T / K-\frac{291.78}{T / K}$
$l_{12}=-25.572+6.5240 \times 10^{-2} T / K-5.4600 \times 10^{-5}(T / K)^{2}+\frac{3322.1}{T / K}$


Fig. 3. Isobaric VLE for ethanol (1) - heptane (2) at 101.33 kPa ; $(\bigcirc)$ ): This work; $(\mathbf{\Delta}, \triangle)$ : Katz and Newman [6], 101 kPa ; $(\square, \square)$ : Raal et al. [7], $101 \mathrm{kPa} ;(-)$ ) : PRSV equation.


Fig. 4. Isothermal VLE for ethanol (1) - heptane (2) at (343 to 483) K; (©, ○): This work, $363.15,393.15,423.15 \mathrm{~K} ;(\boldsymbol{\nabla}, \nabla)$ : Ramaloh and Delmas [16], 343, $363 \mathrm{~K} ;(\mathbf{\Delta}, \triangle)$ : Berro et al. [12], 343.17 K ; $\square, \square$ ): Seo et al. [17], 483.15 K ; (-): PRSV equation.

Table 8
Parameters in AS mixing rule and coefficients for molar Gibbs energy in Eyring equation.

| Temperature <br> $T / \mathrm{K}$ | Parameters in AS mixing rule <br> $k_{12}\left(=k_{21}\right) /-$ |  | Coefficient for molar <br> $l_{12}\left(=-l_{21}\right) /-$ |
| :--- | :---: | :---: | :---: |
| 343.17 | 0.087 | 0.059 |  |
| 363.15 | 0.108 | 0.069 | 0.08 |
| 393.15 | 0.128 | 0.072 | -0.09 |
| 423.15 | 0.165 | 0.105 | -0.28 |
| 483.15 | 0.172 | 0.068 |  |



Fig. 5. RDs of VLE from calculation by PRSV equation for ethanol (1) - heptane (2); (○): $363.15 \mathrm{~K} ;(\bigcirc): 393.15 \mathrm{~K}$; ( $\square$ ): 423.15 K ; (+): 101.33 kPa .

Using Eqs. (12)-(13), isobaric VLE were calculated at 101.33 kPa (Fig. 3). Tables 5 and 6 provides a summary of calculated results. Fig. 5 shows the RDs for temperature and vapor phase mole fraction. The AARDs (Table 5) were $0.282 \%$ and $2.741 \%$ for the temperature and the mole fraction in the vapor phase, respectively. Though the reproducibility was poor at low mole fractions in vapor phase, the calculations agreed with the experimental data including the azeotrope.

### 4.6. Correlation of viscosity

Brunson and Byers have reported the liquid viscosity for ethanol (1) - heptane (2) with dynamic laser light scattering technique [29], that seems to be at conditions close to the saturation pressure. Measurements were carried out at mole fractions of ethanol of $x_{1}=0.0,0.3,0.5,0.8,0.9$ and 1.0. In this research, values of viscosity at $363.15 \mathrm{~K}, 393.15 \mathrm{~K}$ and 423.15 K , were obtained by interpolation of literature data. In the interpolation, the functional form of the viscosity at a given mole fraction, was assumed to be cubic in inverse temperature:
$\ln \eta_{e s t}\left(T, p, x_{1}\right)=a_{v i s 0, x_{1}}+a_{v i s 1, x_{1}}\left(\frac{1}{T}\right)+a_{v i s 2, x_{1}}\left(\frac{1}{T}\right)^{2}+a_{v i s 3, x_{1}}\left(\frac{1}{T}\right)^{3}$

Viscosity data were calculated by using the PRSV equation combined with Eyring theory. In Eq. (7), the viscosity of pure ethanol or heptane, $\eta_{e s t, i}{ }^{\circ}(T, p)$ at a given pressure, $p$, was not available. Therefore, the following assumptions were employed:

$$
\begin{equation*}
\eta_{e s t, i}^{\circ}(T, p)=\eta_{e s t}\left(T, p_{i}^{s}, x_{i}=0\right) \quad(i=1,2) \tag{15}
\end{equation*}
$$

Table 9
Estimated and calculated liquid phase viscosities for ethanol (1) - heptane (2) at isothermal conditions.

| Mole fraction of ethanol $x_{1} /-$ | ${ }^{a}$ Estimated viscosity <br> $\eta_{\text {est }} / \mathrm{mPa} \cdot \mathrm{s}$ | Calculated $\sigma=0.05 /-$ <br> $\eta_{\text {cal }} / \mathrm{mPa}$ s | cosity ${ }^{\mathrm{b}} \Delta \eta / \mathrm{mPa} \cdot \mathrm{~s}$ | ${ }^{\mathrm{c}}\|\Delta \eta\| / \eta_{\text {est }} / \%$ | T-dependent $\eta_{\text {cal }} / \mathrm{mPa} \cdot \mathrm{s}$ | ${ }^{\text {b }} \Delta \eta / \mathrm{mPa} \cdot \mathrm{s}$ | ${ }^{\mathrm{c}}\|\Delta \eta\| / \eta_{\text {est }} / \%$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 363.15 K |  |  |  |  |  |  |  |
| 0.0 | 0.198 |  |  |  |  |  |  |
| 0.3 | 0.221 | 0.232 | 0.011 | 5.081 | 0.220 | -0.001 | 0.295 |
| 0.5 | 0.254 | 0.266 | 0.012 | 4.561 | 0.249 | -0.005 | 1.902 |
| 0.8 | 0.337 | 0.346 | 0.009 | 2.668 | 0.332 | -0.005 | 1.412 |
| 0.9 | 0.375 | 0.386 | 0.011 | 3.018 | 0.377 | 0.002 | 0.683 |
| 1.0 | 0.438 |  |  |  |  |  |  |
|  |  | average | 0.011 | 3.832 | average | -0.003 | 1.073 |
| 393.15 K |  |  |  |  |  |  |  |
| 0.0 | 0.160 |  |  |  |  |  |  |
| 0.3 | 0.167 | 0.173 | 0.006 | 3.778 | 0.176 | 0.009 | 5.043 |
| 0.5 | 0.193 | 0.189 | -0.004 | 2.313 | 0.196 | -0.003 | 0.564 |
| 0.8 | 0.239 | 0.231 | -0.008 | 3.174 | 0.234 | -0.005 | 1.869 |
| 0.9 | 0.257 | 0.254 | -0.003 | 1.046 | 0.256 | -0.001 | 0.228 |
| 1.0 | 0.284 |  |  |  |  |  |  |
|  |  | average | -0.002 | 2.578 | average | 0.000 | 1.926 |
| 423.15 K |  |  |  |  |  |  |  |
| 0.0 | 0.134 |  |  |  |  |  |  |
| 0.3 | 0.130 | 0.130 | 0.000 | 0.312 | 0.141 | 0.011 | 8.879 |
| 0.5 | 0.141 | 0.134 | -0.007 | 4.937 | 0.147 | 0.006 | 3.697 |
| 0.8 | 0.174 | 0.156 | -0.018 | 10.283 | 0.165 | -0.009 | 5.317 |
|  | 0.182 | 0.169 | -0.013 | 6.875 | 0.175 | -0.007 | 3.838 |
| $1.0$ | 0.187 |  |  |  |  |  |  |
|  |  | average | -0.009 | 5.602 | average | -0.001 | 5.433 |

> a interpolated data of Bernson and Byers [29].
> b $|\Delta \eta| / y_{1, \text { exp }}=\frac{\left|\eta_{c a t}-\eta_{\text {est }}\right|}{\eta_{\text {est }}} \times 100$
> c $\Delta \eta=\eta_{\text {cal }}-\eta_{\text {est }}$


Fig. 6. Liquid viscosity for ethanol (1) - heptane (2);(O, ©, $\square$ ): Brunson and Byers [29], 363.15, 393.15, 423.15 K (interpolated); (——): PRSV equation with $\sigma=0.05$; (—): PRSV equation with $\sigma=5.7055-1.0384 \times 10^{-2} T / K-\frac{673.49}{T / K}$.
where $\eta_{i}\left(T, p_{i}^{s}, x_{i}=0\right)$ is the viscosity at the saturated vapor pressure of pure ethanol (1) or heptane (2), and available in the literature [29]. In Eq. (8), the parameter $\sigma$ was determined under the temperature independent and dependent assumptions. The temperature independent parameter was $\sigma=-0.05$. The parameters, separately fitted with the viscosity data at $363.15 \mathrm{~K}, 393.15 \mathrm{~K}$ and 423.15 K , were listed in Table 8. Calculated viscosity values are shown in Fig. 6 and Table 9. The temper-


Fig. 7. RDs of viscosity from calculation by PRVV equation for ethanol (1) heptane (2); ( $\bigcirc$ ): $363.15 \mathrm{~K}, \sigma=0.08$; ( $\square$ ): $393.15 \mathrm{~K}, \sigma=-0.09$; ( $\triangle$ ): 423.15 K , $\sigma=-0.28$; ( ) : $363.15 \mathrm{~K}, \sigma=0.05$; ( $\boldsymbol{\square}): 393.15 \mathrm{~K}, \sigma=0.05$; ( $\mathbf{\wedge}): 423.15 \mathrm{~K}, \sigma=0.05$.
ature dependent parameters provided better agreements with the liquid viscosity data than those with the temperature independent parameter especially in low temperature range. The temperature dependence of $\sigma$ was given by:
$\sigma=5.7055-1.0384 \times 10^{-2} T / K-\frac{673.49}{T / K}$
Though some deviations were found in the calculation at 423.15 K , it seems to be related to system oscillations. The AARDs (Fig. 7) with $\sigma=-0.05$ were $3.832 \%, 2.578 \%$ and $5.602 \%$ at $363.15 \mathrm{~K}, 393.15 \mathrm{~K}$ and 423.15 K , respectively. Otherwise, those with Eq. (16) were $1.073 \%$, $1.926 \%$ and $5.433 \%$ at $363.15 \mathrm{~K}, 393.15 \mathrm{~K}$ and 423.15 K , respectively.

## 5. Conclusions

Vapor-liquid equilibria data were measured for binary mixtures of ethanol - heptane with a modified Rose-Williams equilibrium still at one isobar and at three isotherms, ( 363 to 423 ) K, corresponding to conditions for direct fuel injection into gasoline engines. The PRSV equation
and the Adachi-Sugie mixing rules provided correlation of the data except at ethanol mole fractions lower than 0.25 . Combination of Eyring theory with an additional parameter, $\sigma$, allowed estimation of liquid viscosity data at ( 363 to 423 ) K. The data and correlation are useful for the design of direct fuel injectors that use bioethanol blended gasolines. Considering corresponding state correlations for transport properties, the methods outlined in this paper can probably be extended to the correlation of thermal conductivities and diffusion coefficients of bioethanol-gasoline mixtures.

## Credit authorship contribution statement

Tomoya Tsuji: Calculation, Collection and analysis of literature, Data analysis, Writing - Review \& Editing. Hiroyuki Matsuda: Methodology, Investigation, Writing - Review \& Editing. Makoto Kawakita: Investigation, Writing- Original Draft. Kiyofumi Kurihara: Design of equilibrium still and ebulliometer, Co-supervision of MK. Katsumi Tochigi: Overall Planning, Data analysis, Supervision of MK.

## Declaration of Competing Interest

The authors declare that they have no conflicts of interest with the contents of this article.

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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## Supplementary materials

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